

# Quantum Coulomb Glass on the Bethe Lattice



Technische Universität München



Munich Center of Quantum  
Science and Technology

Izabella Lovas

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Collaborators: Annamária Kiss, Catalin Pascu Moca, Gergely Zaránd

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# Localization transition of fermions in the presence of disorder

Interacting systems with long-range Coulomb interactions: **Coulomb Glass**

## Peculiar properties of interacting systems

Reduced screening ➡ Anomalies already on *metallic* side

Depleted density of states at Fermi energy on *insulating* side dimension

Obeying Efros-Shklovskii scaling  $\rho(\varepsilon) \sim \varepsilon^{d-1}$

## Coulomb Gap

Glassy dynamics, memory effects...

Efros & Shklovskii (1975), Amir et al. (2011)...

## Classical limit of interacting systems: **Spin glasses**

Insight from (numerically) exactly solvable mean field model ➡ **Sherrington-Kirkpatrick model**  $H = - \sum_{i<j} J_{ij} S_i S_j$

➡ Low T glassy phase, freezed spins: Edwards – Anderson order parameter  $q_{EA} = \overline{m_i^2}$   $m_i = \langle S_i \rangle_H$  Gaussian random variables

➡ Analogue of Coulomb gap: ‘pseudogap’ in the distribution local effective magnetic field  $h_i = \sum_{j:j \neq i} J_{ij} m_j$  ➡  $P(h) \sim h$

➡ Glass melted by thermal fluctuations

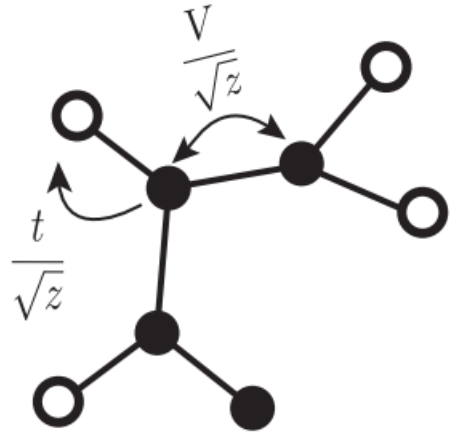
**Quantum Coulomb Glass?**  
**Role of Quantum Fluctuations?**

# Mean Field Electron Glass Model

Disordered interacting fermionic system on a Bethe lattice with coordination #  $z \rightarrow \infty$

$$\hat{H} = -\frac{t}{\sqrt{z}} \sum_{\langle i,j \rangle} (\hat{c}_i^\dagger \hat{c}_j + \text{h.c.}) + \frac{V}{\sqrt{z}} \delta \hat{n}_i \delta \hat{n}_j + \sum_i \varepsilon_i \delta \hat{n}_i$$

coordination number  $\rightarrow$   $\sqrt{z}$   
hopping  $\rightarrow$   $(\hat{c}_i^\dagger \hat{c}_j + \text{h.c.})$   
 $\hat{c}_i^\dagger \hat{c}_i - 1/2$   $\rightarrow$   $\delta \hat{n}_i$   
nearest neighbour repulsion  $\rightarrow$   $\delta \hat{n}_i \delta \hat{n}_j$   
random on-site potential  $\rightarrow$   $\varepsilon_i$



Random site energies with Gaussian distribution  $P(\varepsilon) \sim e^{-\varepsilon^2/(2W^2)}$

A. A. Pastor and V. Dobrosavljević, *Phys. Rev. Lett.* **83**, 4642 (1999).

## Emerging Glassy Phase at Low T

‘Structural’ transition governed by interactions

← Even in the absence of disorder  $W=0$

Exact mapping to a single site DMFT becomes exact

Mean Field Model for Coulomb Glass

↔ Sherrington-Kirkpatrick mean-field model for spin glasses

# Quantum Solution

Most previous studies: - Classical spin glass limit ( $t = 0$ )  
 - Vicinity of phase transition (Landau theory)

Our Work: Quantum Solution Deep in Glassy Phase

## Numerical Ingredients:

- Disorder averaging with replica trick
- Exact mapping to single site problem
- Effective local action:
  - a.) continuous time quantum Monte Carlo (CTQMC)  
 exact but expensive: only in high-T phase
  - b.) iterated perturbation theory (IPT)  
 access to glassy phase

Cavity approach:

sum over nn sites of 0

$$S = \int_{\tau} \bar{c}_{0\tau} (\partial_{\tau} + \epsilon_0) c_{0\tau} + \frac{V}{\sqrt{z}} \sum'_i \int_{\tau} \delta n_{0,\tau} \delta n_{i,\tau} - \frac{t}{\sqrt{z}} \sum'_i \int_{\tau} (\bar{c}_{0\tau} c_{i\tau} + h.c.) + S_{i \neq 0}$$

Integrate out all sites except 0

$$\begin{aligned} S_0^{\text{eff}} = & \int_{\tau} \bar{c}_{0\tau} (\partial_{\tau} + \epsilon_0) c_{0\tau} - t^2 \int_{\tau} \int_{\tau'} \bar{c}_{0\tau} c_{0\tau'} \frac{1}{z} \sum'_i \langle c_{i\tau} \bar{c}_{i\tau'} \rangle_{\text{cav}} + \frac{V}{\sqrt{z}} \sum'_i \int_{\tau} \delta n_{0,\tau} \langle \delta n_{i,\tau} \rangle_{\text{cav}} \\ & - \frac{V^2}{2} \int_{\tau} \int_{\tau'} \delta n_{0,\tau} \delta n_{0,\tau'} \frac{1}{z} \sum'_i \langle \delta n_{i,\tau} \delta n_{i,\tau'} \rangle_{\text{cav}} + \dots \end{aligned}$$

cavity averages in the absence of 0

$$S_0^{\text{eff}} = \int_{\tau} \int_{\tau'} \left\{ \bar{c}_{0\tau} (\delta(\tau - \tau') (\partial_{\tau} + \tilde{\epsilon}_0) - t^2 G(\tau - \tau')) c_{0\tau'} - \frac{V^2}{2} \chi(\tau - \tau') \delta n_{0,\tau} \delta n_{0,\tau'} \right\}$$

Gaussian random field includes Hartree shift

average Green's function

average density correlator

Ensemble of actions + field distribution  $\tilde{P}(\tilde{\epsilon})$ : self-consistent solution

Low T: many symmetry broken states, have to include all of them  
 cavity approach becomes more complicated

→ Replica trick  $\overline{F} \sim \overline{\log Z} = \lim_{n \rightarrow 0} \frac{\overline{Z^n} - 1}{n}$

**We can still map to problem to an ensemble of effective local actions**

$$S_{\text{loc}}(\tilde{\epsilon}) = \int_{\tau} \int_{\tau'} \left\{ \bar{c}_{\tau} \left[ \delta_{\tau, \tau'} [\partial_{\tau'} + \tilde{\epsilon}] - t^2 G(\tau - \tau') \right] c_{\tau'} - \frac{V^2}{2} \tilde{\chi}(\tau - \tau') \delta n_{\tau} \delta n_{\tau'} \right\} - \frac{\beta \tilde{\epsilon}}{2} + \tilde{P}(\tilde{\epsilon})$$

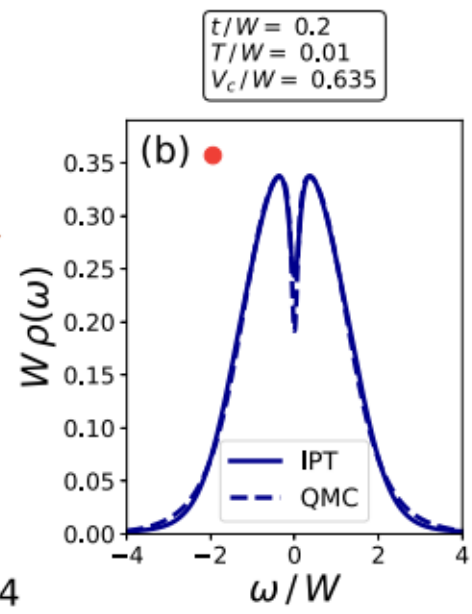
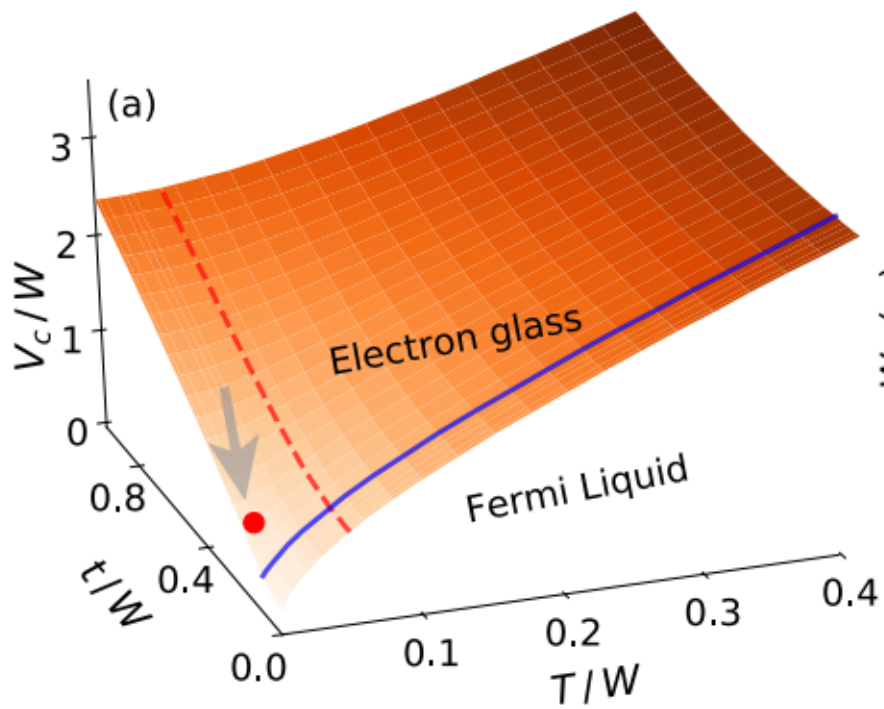
Self-consistency conditions:  $\begin{Bmatrix} G(\tau) \\ \tilde{\chi}(\tau) \end{Bmatrix} = \int d\tilde{\epsilon} \tilde{P}(\tilde{\epsilon}) \begin{Bmatrix} G_{\tilde{\epsilon}}(\tau) \\ \tilde{\chi}_{\tilde{\epsilon}}(\tau) \end{Bmatrix}$

←  $\langle c_{\tau} \bar{c}_{\tau'} \rangle_{\text{loc}}$   
 ←  $\langle \delta n_{\tau} \delta n_{\tau'} \rangle_{\text{loc}} - \langle \delta n_{\tau} \rangle_{\text{loc}}^2$

both calculated from  $S_{\text{loc}}(\tilde{\epsilon})$

**Electron glass transition**

- $\tilde{P}(\tilde{\epsilon})$ : - **Gaussian** in high T phase  
 Variance renormalized by interactions
- **Pseudo-gap** behaviour in glassy phase



*Glass melted by thermal or quantum fluctuations*

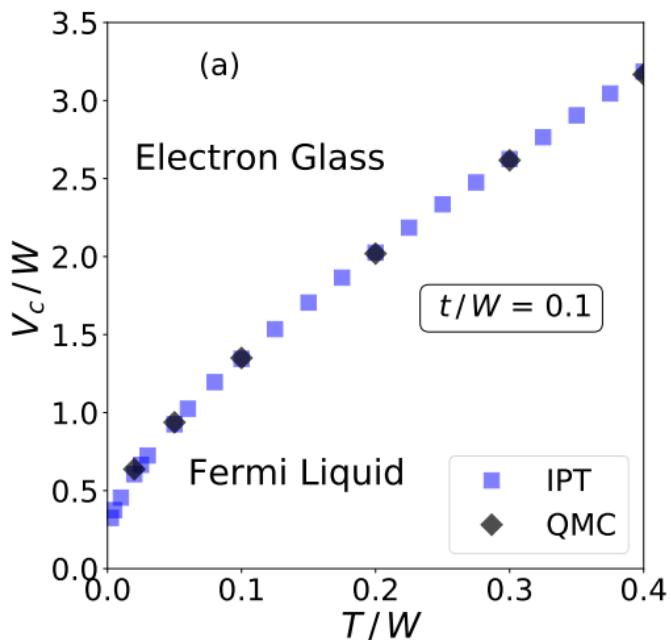
Phase diagram: - continuous time quantum Monte Carlo (CTQMC)  
 - iterated perturbation theory (IPT)

### CTQMC

Expand partition function in hybridization  $t^2 G(\tau - \tau')$

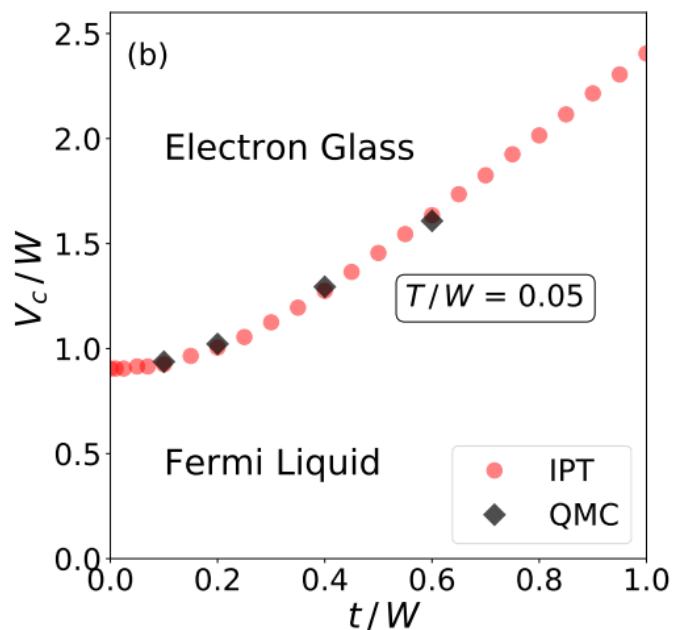
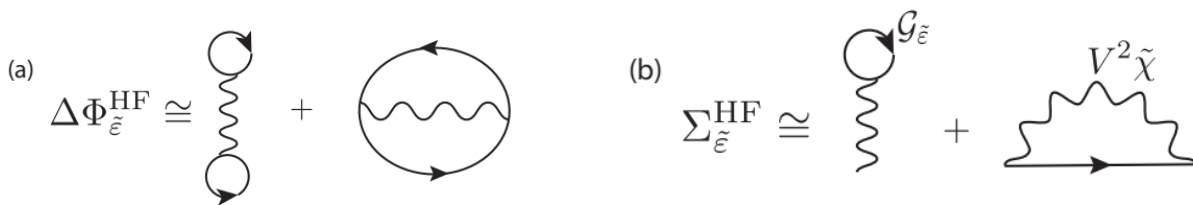
➔ Set of diagrams for observables

Stochastic sampling of diagrams



### IPT

First order expansion in  $V^2 \tilde{\chi}(\tau - \tau') \equiv V^2 (\chi(\tau - \tau') - \langle \delta n \rangle^2)$



# Glassy phase

(IPT results)

# Pseudo-gap opening

Glass Transition Driven by  
Thermal or Quantum Fluctuations

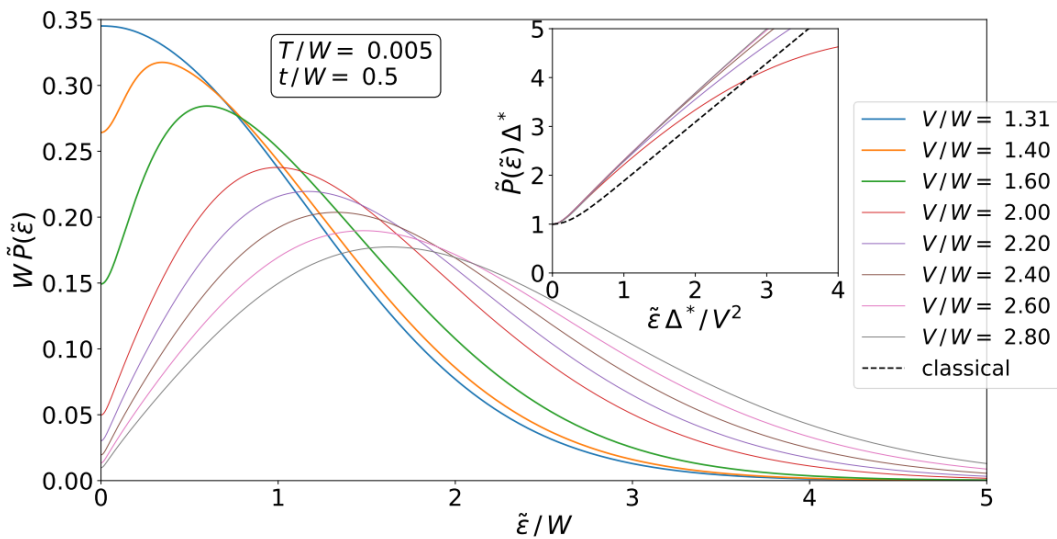
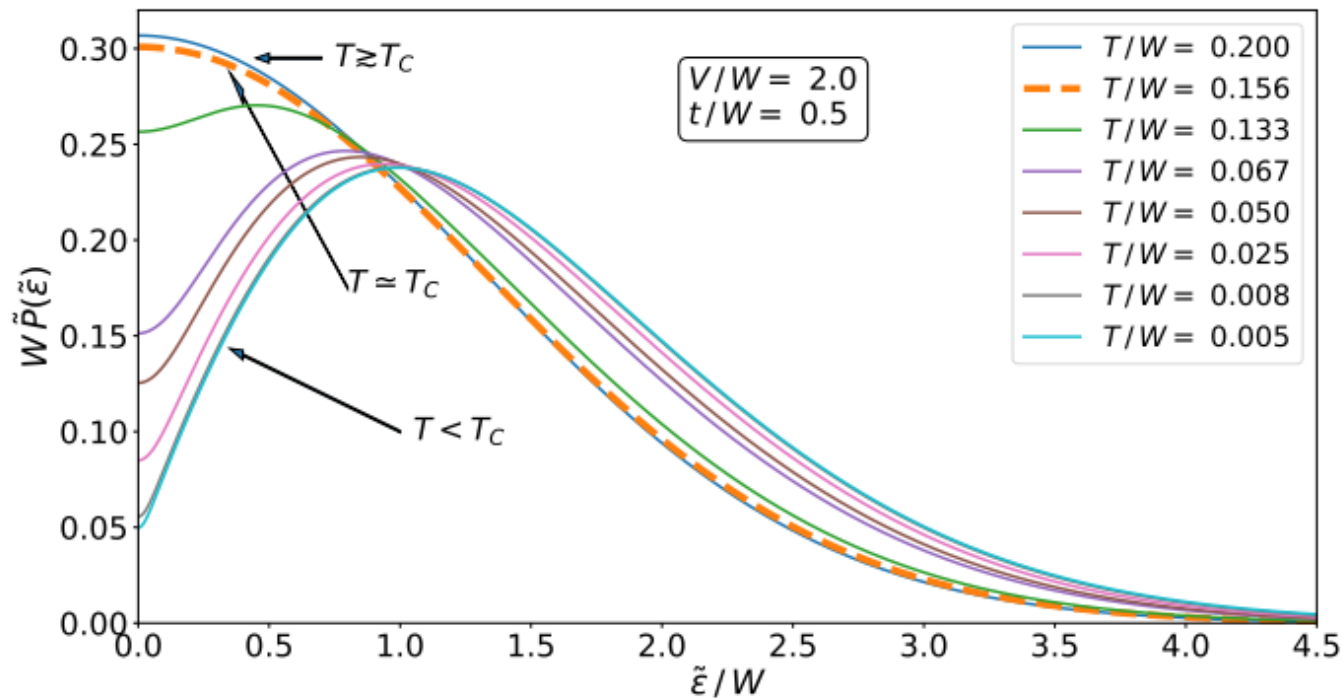
Efros-Shklovskii Scaling

Deep in Glassy Phase

$$\tilde{P}(\tilde{\epsilon}) \approx 1.13 |\tilde{\epsilon}| / V^2$$



Sherrington-Kirkpatrick  
spin glass result



*Universal regime determined by interactions*



*Structural transition driven by interactions alone*



# Pseudo-Gap in Spectral Function

$$\rho(\omega) = \frac{1}{\pi} \text{Im}G(\omega + i0^+) \quad \longleftrightarrow$$

Average Local Density of States  
Accessible in Tunneling Experiments

Efros-Shklovskii Scaling

$$\rho(\omega) \propto \omega/V^2$$

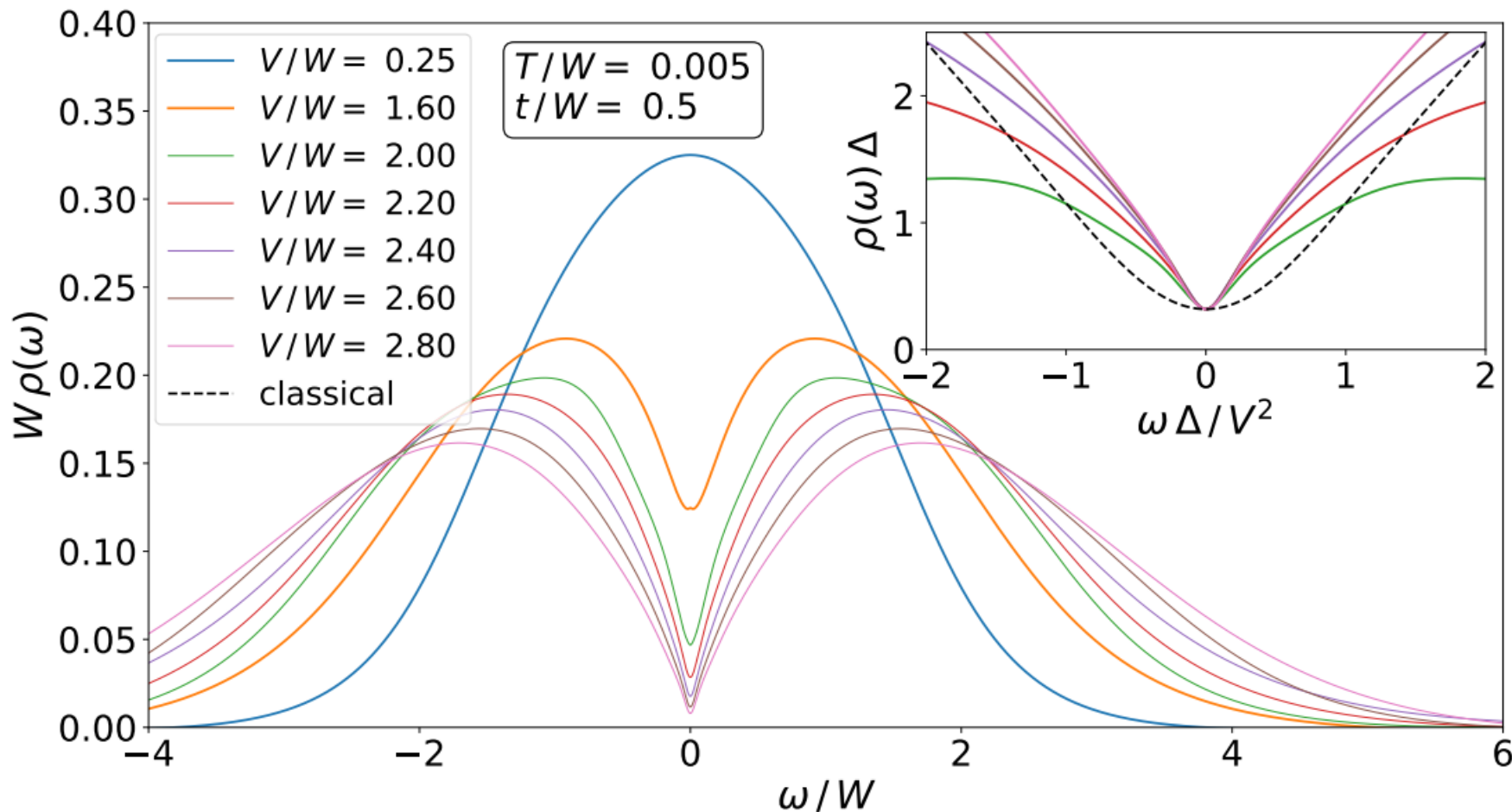
Quantum Scaling Function

$\neq$

Classical Scaling Function



Slight Difference Between  
Quantum and Thermal  
Fluctuations



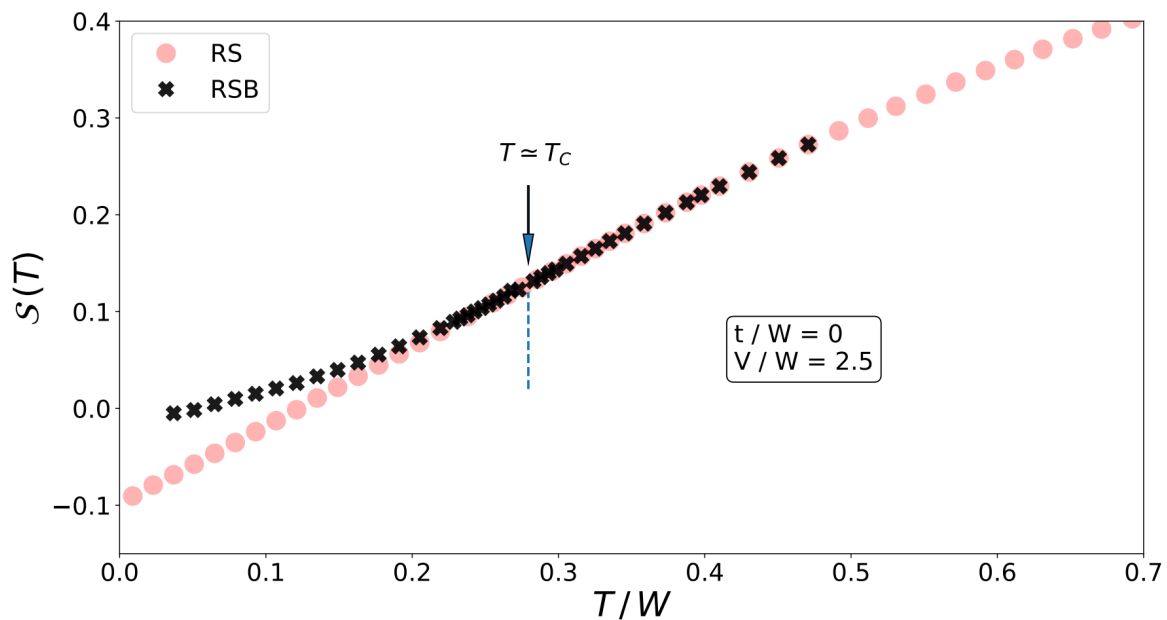
# Thermodynamics

Order of phase transition?  $\mathcal{S} = \frac{\partial \Phi_{\text{latt}}}{\partial T}$

High precision required, not very reliable because of perturbative approach + numerical errors

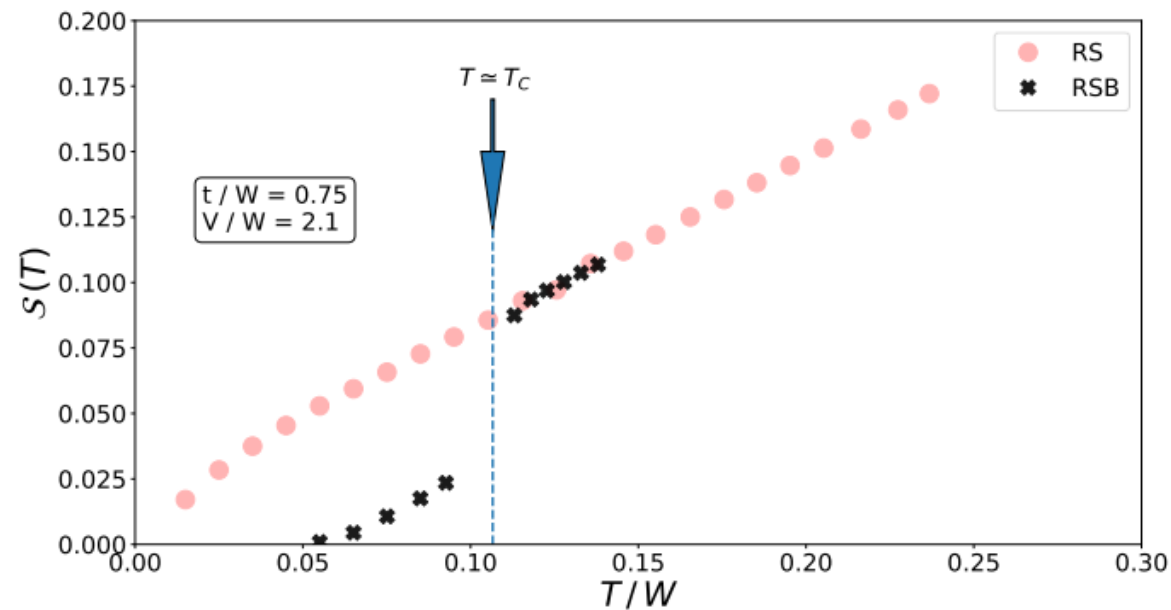
Classical case: (at least) second order transition

Fermi liquid solution unphysical at low T: negative entropy



Quantum case: First order transition

Problem with previous calculations based on Landau approach?



## Summary

### Investigation of a mean field electron glass model in the quantum regime

Phase diagram

Properties of the glassy phase (full replica symmetry breaking)

- Pseudogap in spectral function and in distribution of effective Hartree fields
- Universal regime: Efros-Shklovkii scaling
- First order transition away from classical limit (?)

*Example for structural transition, governed by interactions (not disorder)*

## Outlook / Open questions

- Numerically exact solution with CTQMC in glassy phase: work in progress
- Role of spin degrees of freedom
- Coulomb glass in 2D/3D?
- Dynamics, anomalously slow response

Thank you for your attention!